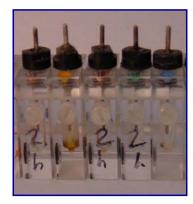
Simulation of Membrane Potentials in Ion Selective Electrodes

Title Slide

Mapping Slide Are Magnus Bruaset Simula Research Laboratory Oslo, Norway

Tomasz Sokalski Åbo Academy / ProSense Åbo (Turku), Finland

June 18, 2004



[ProSense]

This talk outlines a solver for coupled Nernst-Planck and Poisson equations



Ion selective electrodes (ISEs)



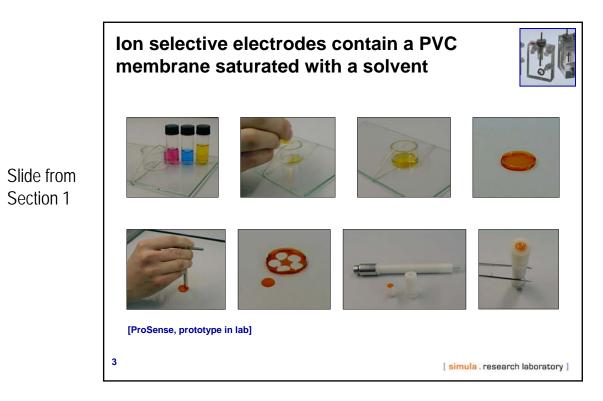
Mathematical model:
Nernst-Planck and Poisson

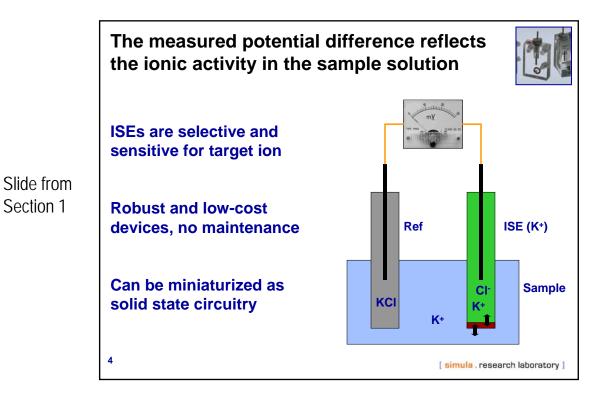


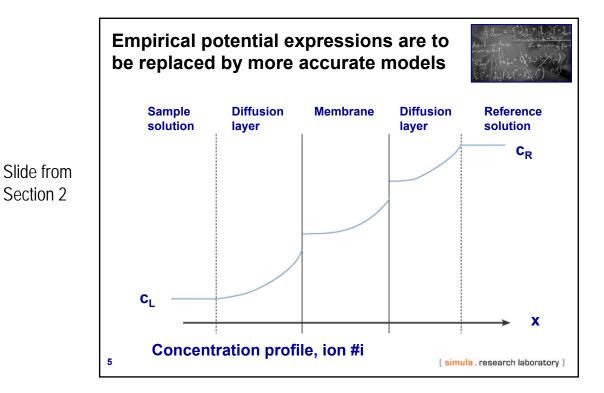
Simulator design

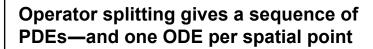
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$$\frac{\partial c_i(x,t)}{\partial t} = -\frac{\partial}{\partial x} [f_i(x,t)] \quad \text{for } x \in \Omega$$

$$f_i(x,t) = -D_i(x,t) \frac{\partial c_i(x,t)}{\partial x} - K_i(x,t)c_i(x,t),$$

$$K_i(x,t) = -D_i(x,t)z_i \frac{F}{RT} E(x,t)$$

$$arepsilon rac{\partial E(x,t)}{\partial t} = I - F \sum_{i=1}^{n_{ ext{ion}}} z_i f_i(x,t) \quad ext{for } x \in \Omega$$

Concentration c_i of specie #i, i=1,...,n_{ion}

Unknowns c_i and E

Electrical field E

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Operator splitting gives a sequence of PDEs—and one ODE per spatial point



Slide from Section 2 (Animation)

$$\frac{\partial c_i(x,t)}{\partial t} = -\frac{\partial}{\partial x} \left(f_i(x,t) \right) \quad \text{for } x \in \Omega$$

$$f_i(x,t) = -D_i(x,t) \frac{\partial c_i(x,t)}{\partial x} - K_i(x,t)c_i(x,t),$$

$$K_i(x,t) = -D_i(x,t)z_i \frac{F}{RT} E(x,t)$$

$$\varepsilon$$
 $E(x,t) = I - F \sum_{i=1}^{n_{\text{ion}}} z_i f_i(x,t) \quad \text{for } x \in \Omega$

Concentration c_i of specie #i, i=1,...,n_{ion}

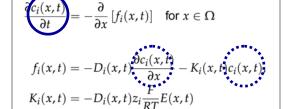
Unknowns c_i and E

Electrical field E

[simula , research laboratory]

Operator eplitting gives a seguence of





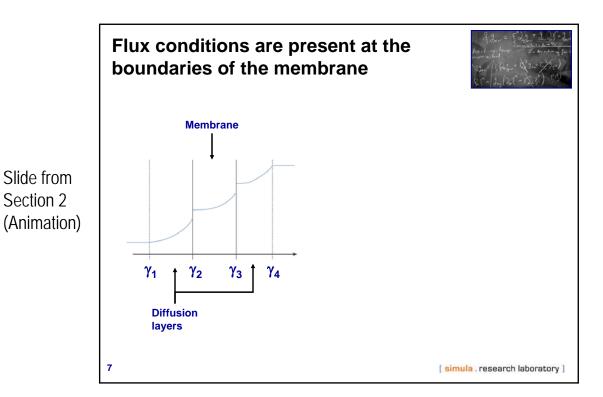
Concentration c_i of specie #i, i=1,...,n_{ion}

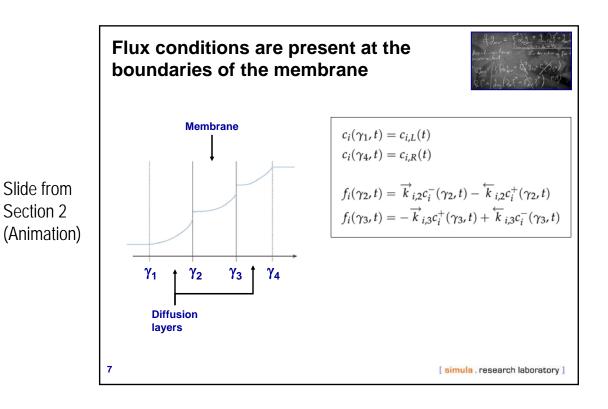
Slide from Section 2 (Animation)

$$\varepsilon \frac{\partial E(x,t)}{\partial t} = I - F \sum_{i=1}^{n_{\mathrm{ion}}} \mathbf{z} f_i(x,t) \quad \text{for } x \in \Omega$$

Electrical field E

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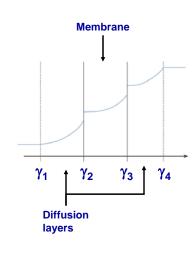




Flux conditions are present at the boundaries of the membrane



Slide from Section 2 (Animation)



$$c_i(\gamma_1, t) = c_{i,L}(t)$$

$$c_i(\gamma_4, t) = c_{i,R}(t)$$

$$f_i(\gamma_2, t) = \overrightarrow{k}_{i,2}c_i^-(\gamma_2, t) - \overleftarrow{k}_{i,2}c_i^+(\gamma_2, t)$$

$$f_i(\gamma_3, t) = -\overrightarrow{k}_{i,3}c_i^+(\gamma_3, t) + \overleftarrow{k}_{i,3}c_i^-(\gamma_3, t)$$

Membrane only

$$f_{i}(\gamma_{2},t) = \overrightarrow{k}_{i,2}c_{i,L}(t) - \overleftarrow{k}_{i,2}c_{i}(\gamma_{2},t)$$

$$f_{i}(\gamma_{3},t) = -\overrightarrow{k}_{i,3}c_{i,R}(t) + \overleftarrow{k}_{i,3}c_{i}(\gamma_{3},t)$$

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In summary, PDE-based models offer more versatile simulations of ionic activity

Conclusion Slide

Computational use of PDEs in this context is novel

Good match between computations and experiments is observed

Use of finite elements allows models to capture ignored effects



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